

# Homework 1 - Solutions

19th October 2004

## 1

Entropy is the logarithm of the number of accessible states. If there are  $N$  chairs and  $k$  students then since students are distinguishable the number of configurations they can be in is  $N!/(N-k)!$  and the entropy  $\sigma = \log(N!/(N-k)!)$ . Taking say  $N = 70$  and  $k = 40$  we have  $\sigma \approx 70 \log 70 - 70 - 30 \log 30 + 30 \approx 155$ . (Don't forget, the log is the natural logarithm.)

## 2

Define a set of configurations of the system by specifying which particles are in which side of the container. Such a configuration changes when a particle crosses the (imaginary) border separating left and right. Let the average time between changes be  $t$ . There is one configuration where all particles are on the left. There are  $M = 10^{N_0/3}$  other configurations. Starting with random configuration, it takes on average approximately  $M$  steps to get to the desired configuration. Thus the time it takes to get there is  $Mt$ . Thus we just need to find  $t$  and multiply it by  $M$ .

Suppose  $L$  is the size of the box. If there were only one molecule it would cross the border (between right and left) approximately in time  $L/v$  where  $v$  is its velocity. If there are  $N_0$  molecules it would happen  $N_0$  times more often so the time is  $t = L/(vN_0)$ .

Now you can plug in the numbers, but from the values given you can clearly see that it is in the range  $(10^{-100}, 10^{100})$  even in the ages of the universe. Multiplying  $t$  and  $M$  we get for the final answer  $T = 10^{2 \times 10^{23} \pm 100} \approx 10^{2 \times 10^{23}}$ .

### 3

$$g(U, N, V) = CU^{\alpha N} V^{\beta N}$$

By definitions

$$\frac{1}{\tau} = \left( \frac{\partial \log g}{\partial U} \right)_{NV} = \alpha N \frac{\partial \log U}{\partial U} = \frac{\alpha N}{U}$$
$$p = \tau \left( \frac{\partial \log g}{\partial V} \right)_{NU} = \frac{\beta U}{\alpha V}$$

About ideal gas we know the following. Its energy is  $U = \frac{3}{2}NkT$  and it satisfies the ideal gas equation  $pV = NkT$ . Heat capacity at constant volume is  $C_v = \left( \frac{\partial U}{\partial T} \right)_{VN} = \frac{3}{2}Nk$ .

In our gas, from the equation for the temperature above we have  $U = \alpha NkT$  and comparing with ideal gas energy this gives  $\alpha = 3/2$ . You can also get this by expressing  $S$  in terms of  $T, V, N$  and then using  $C_v = T(\partial S/\partial T)_{VN}$ .

From the above equation for  $p$  we have  $pV = U\beta/\alpha = \alpha N\tau\beta/\alpha = NkT\beta$  which comparing to the ideal gas equation gives  $\beta = 1$ .

### 4

$$\sigma(s) \approx \sigma_0 - 2s^2/N$$

One spin has moment  $m$  and energy  $mB$ . There are  $2s$  more spins up than down. So together they have moment  $M = 2sm$  and energy  $U = 2smB$ . Substituting  $s = U/2mB$  into above we get  $\sigma(U)$ . Differentiating with  $U$  we get  $1/\tau$ . From  $M = 2sm$ ,  $M/mN = 2s/N = U/mBN$ .

### 5

$$\sigma(n, N) = \log(g(N, n)) = \log(N+n-1)! - \log(N-1)! - \log n! \approx (N+n) \log(N+n) - N \log N - n \log n$$

Substituting  $n = U/\hbar\omega$

$$\sigma(U, N) \approx (N + U/\hbar\omega) \log(N + U/\hbar\omega) - N \log N - (U/\hbar\omega) \log(U/\hbar\omega)$$

$$\frac{1}{\tau} = \frac{\partial \sigma(U, N)}{\partial U} = \dots = \frac{1}{\hbar\omega} \log \frac{N + U/\hbar\omega}{U/\hbar\omega}$$

$$U = \frac{N\hbar\omega}{\exp(\hbar\omega/\tau) - 1}$$

6

$$p_n = \frac{N!}{n!(N-n)!} p^n (1-p)^{N-n} \quad (1)$$

$$= \frac{N(N-1)\cdots(N-n+1)}{n!} p^n (1-p)^{N-n} \quad (2)$$

$$\rightarrow \frac{(pN)^n}{n!} (1-p)^N \quad (3)$$

$$= \frac{r^n}{n!} ((1-p)^{1/p})^{Np} \quad (4)$$

$$\rightarrow \frac{r^n}{n!} e^{-r} \quad (5)$$

Another solution:

We will use the following fact. From Taylor expansion

$$\log(1 - n/N) = -n/N + (n/N)^2/2 + \dots$$

The three dots " $\dots$ " below mean that the term goes to zero in the limit we are considering. We have

$$\begin{aligned} \log \frac{N!}{n!(N-n)!} &= \log N! - \log(N-n)! - \log n! \\ &\approx N \log N - N - (N-n) \log(N-n) + (N-n) - \log n! \\ &= N \log N - N(1 - n/N) \log N(1 - n/N) - n - \log n! \\ &= N \log N - N \log N - N \log(1 - n/N) + n \log N + n \log(1 - n/N) - n - \log n! \\ &= n + \dots + n \log N + \dots - n - \log n! \\ &\approx n \log N - \log n! \end{aligned}$$

$$\begin{aligned} \log p^n (1-p)^{N-n} &= n \log p + (N-n) \log(1-p) \\ &\approx n \log p - Np + np \\ &\approx n \log p - r \end{aligned}$$

Summing the two log's we have

$$\log p_n = n \log(Np) - r - \log n!$$

$$p_n = \frac{r^n}{n!} e^{-r}$$

Let's calculate the mean and the variance. From Taylor expansion

$$e^r = \sum_{n=0}^{\infty} \frac{r^n}{n!}$$

Using this in  $\sum p_n$  we get  $\sum p_n = 1$  as we should for a probability distribution (If we didn't we would divide it by the result of this summation).

The mean is  $\langle n \rangle = \sum np_n$

The variance is  $\langle (\Delta n)^2 \rangle = \langle (n - \langle n \rangle)^2 \rangle = \langle n^2 - 2n \langle n \rangle + \langle n \rangle^2 \rangle = \langle n^2 \rangle - 2 \langle n \rangle^2 + \langle n \rangle^2 = \langle n^2 \rangle - \langle n \rangle^2$

And  $\langle n^2 \rangle = \sum n^2 p_n$ .

To calculate these we do the following. In the expansion of  $e^r$  above, take  $r(d/dr)$  of both sides. On the left we get  $re^r$ , on the right  $\sum nr^n/n! = e^r \sum np_n = e^r \langle n \rangle$ . Thus  $\langle n \rangle = r$ .

Similarly take  $r(d/dr)(r(d/dr))e^r$ . It gives on the left side  $(r^2 + r)e^r$  and the right side  $\sum n^2 r^n/n! = e^r \langle n^2 \rangle$  so  $\langle n^2 \rangle = r^2 + r$ .

Thus  $\langle (\Delta n)^2 \rangle = \langle n^2 \rangle - \langle n \rangle^2 = r^2 + r - r^2 = r$ .

Fractional uncertainty is  $\sqrt{\langle (\Delta n)^2 \rangle} / \langle n \rangle = 1/\sqrt{r}$ .